PHYSICA D

Successive homoclinic tangencies to a limit cycle

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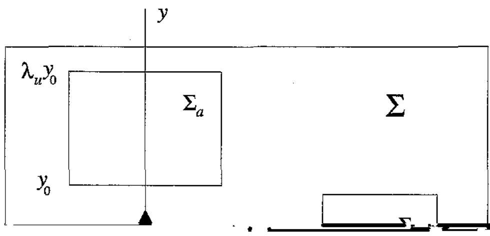
Abstract

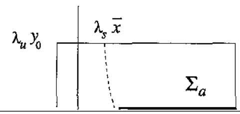
The dynamics near a perturbed degenerate homoclinic connection to a periodic orbit in three dimensions is modeled by



2.1. Boundary conditions

An important feature of the definition of the return sections Σ_a and Σ_b is that the linear map L takes the lower boundary $y = y_0$ of Σ_a to its upper boundary $y = \lambda_u y_0$, while L takes the right boundary $x = x_0$ of Σ_b onto its left boundary $x = \lambda_s x_0$. Thus a point $A =$

and minimum on the interval $[y_0, \lambda_u y_0]$, as for example,



4. Simple fixed points

The simplest periodic orbits of the flow are the trajectories which link up with themselves after only one pass through the global region of the flow. These correspond to fixed points $(x, y) \in \Sigma_a$ of our mapping which satisfy $(x, y) = L^m G(x, y)$, or more explicitly

$$x = \lambda_s^m \phi(y)\,,$$
$$y = \lambda_u^m \left\{ \left[\mu + \gamma \bar{x}(x, y)\right] y/y_0 + \epsilon f(y) \right\}. \qquad (20)$$

The substitution of the first equation $x = \lambda_s^m \phi(y)$ into the second equation yields an equation for y decoupled

fixed points merge again in a second saddle-node tangency and disappear. Thus as μ varies from positive to negative, two cycles are created in a saddle-node bifurcation and then the same two cycles merge and are destroyed in another saddle-node bifurcation.

Now we find where the saddle-node bifurcations occur relative to the primary homoclinic tangencies. The saddle nodes occur when

$$\beta = -\epsilon f'(y)\,, \qquad (24)$$

and a simple computation shows that the corresponding fixed point $(\lambda_s^m \phi(y), y)$ is also given by a solu-

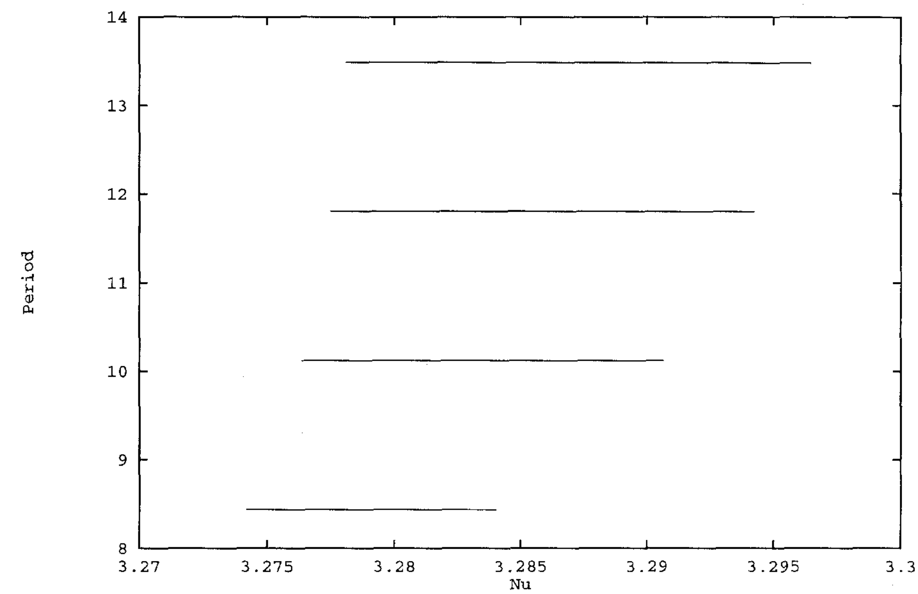


Fig. 10. Plot of period versus ν at $\eta = -5.5$ for periodic orbits corresponding to tongues with rotation numbers 1/5, 1/6, 1/7 and 1/8, [redacted]

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